

Structural Variation of Silver Clusters from Ag₁₃ to Ag₁₆₀

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The structures of silver clusters from Ag₁₂₁ to Ag₁₆₀ were optimized with a modified dynamic lattice searching (DLS) method, named as DLS with constructed core (DLSc). The interaction among silver atoms is modeled by the Gupta potential. Structural characteristic of silver clusters with the growth of cluster size is investigated with the newly optimized structures and our previous results from Ag₁₃ to Ag₁₂₀. A set of amorphous structures was obtained in the size range of 13–48, together with several ordered structures. The putative stable motif is an icosahedron from Ag₄₉ to Ag₆₁ and then changes to a decahedron in the size range of 62–160. Some of the results are consistent with experiments. Furthermore, it was also found that, for clusters with decahedral motif, the stable structure is a result of the competition among the different Marks decahedral motifs. On the other hand, different from the Lennard-Jones cluster, there are some silver clusters with the face-centered cubic (fcc) motif in the size range of 13–160. But the fcc motif can only be obtained for some specific sizes.

1. Introduction

There have been great interests in experimental and theoretical investigations of small silver cluster due to its applications in photography, catalysis, electronic materials, surface nanostructuring, etc.^{1–6} The study of the structural changes of small (tens through hundreds of atoms) silver clusters with their sizes is a hotpot because it may reveal the relationship between structures and their physical and chemical properties.^{7,8}

In experiments, the observed structure of a silver cluster is a result of the competition between thermodynamic and kinetic factors.^{9–12} At small size, silver clusters take the forms of icosahedral (Ih), truncated Marks decahedral (*m*-Dh), face-centered cubic (fcc), and amorphous motifs. Many experiments and theoretical studies have been devoted to determination of the putative stable structures depending on the cluster size. In 1991, Hall et al.⁹ studied the structures of small silver clusters grown in inert gas aggregation (IGA)¹³ source using electron diffraction. Ih, Dh, and fcc clusters between 2 and 4 nm diameter were observed. It was found that the relative proportions of these motifs depend on experimental conditions, and for small clusters (~2 nm of diameter, corresponding to 100–200 atoms), the structures are mainly in Dh motif. Then, in 1997, Reinhard et al.¹⁰ studied the size-independent fcc-to-icosahedron structural transition of large silver clusters grown in IGA source up to 11 nm in diameter. On the basis of these works, Baletto et al.^{11,14,15} did a series of works on theoretical simulation of the silver clusters. In 2000, they studied the growth of free silver clusters up to 150 atoms by molecular dynamics simulation from a small seed.¹¹ It was shown that Dh is the most possible motif in the most experimental conditions, but Ih structures are indeed energetically favorable. In 2001, they investigated the growth of silver clusters up to 600 atoms by molecular dynamics simulation on realistic time scales and in a temperature range from 400 to 650 K.¹⁴ At low and intermediate temperatures (350 < *T* < 500 K), the clusters grow through a sequence from Ih

to Dh and then back to Ih. The first Ih to Dh transition takes place always in the size range of 55–75, while the crucial transformation from Dh to Ih takes place around Ag₁₀₀ if *T* < 400 K and around Ag₂₀₀ if *T* > 400 K. At high temperatures (*T* > 600 K), most of the clusters take the structural motif of fcc, after passing through a Dh regime around Ag₂₀₀. Then, they studied the crossover size among structural motifs, which showed that the best motif is Ih for the clusters size smaller than 147, Dh in the size range of 300–20 000, and fcc at much larger size.¹⁵ Finally, they explained the silver cluster structures by energy, thermodynamics and kinetics.¹² These simulations provided a great help for a deep understanding of the experimental observations. Moreover, other works on theoretical analysis of silver clusters have also been reported.^{16–21}

On the other hand, optimization algorithms, such as the evolutionary algorithm,²² random tunneling algorithm (RTA),²³ and dynamic lattice searching (DLS) methods²⁴ have been developed to study the putative stable structures and the corresponding lowest potential energies of silver clusters at given sizes. However, limited by the efficiency of the algorithms, only small silver clusters or the clusters with given specific sizes were optimized, e.g., in the work of Baletto et al.,¹⁵ only the energies of the silver clusters at the magic numbers were investigated. The optimization of silver clusters in a consecutive size range is still scarce. In our previous works,^{23,24} the putative stable structures of silver clusters up to 80 atoms and from 61 to 120 atoms were studied, respectively. It was found that the dominating motifs are disordered morphologies in the size range of 15–47, Ih in the size range of 48–54, and the dominating motif changes from Ih to Dh in the size range of 55–75. It was also found that, from Ag₆₁ to Ag₁₂₀, the majority of the clusters have Dh motif, and these results are found to be consistent with experimental results of Hall,⁹ but different to the theoretical results of Baletto.¹⁵

In this work, a modified DLS^{25,26} method, named as DLS with constructed core (DLSc),²⁷ was used to optimize the structures of the silver clusters from Ag₁₂₁ to Ag₁₆₀. It was found that the majority of the structures in this size range are also

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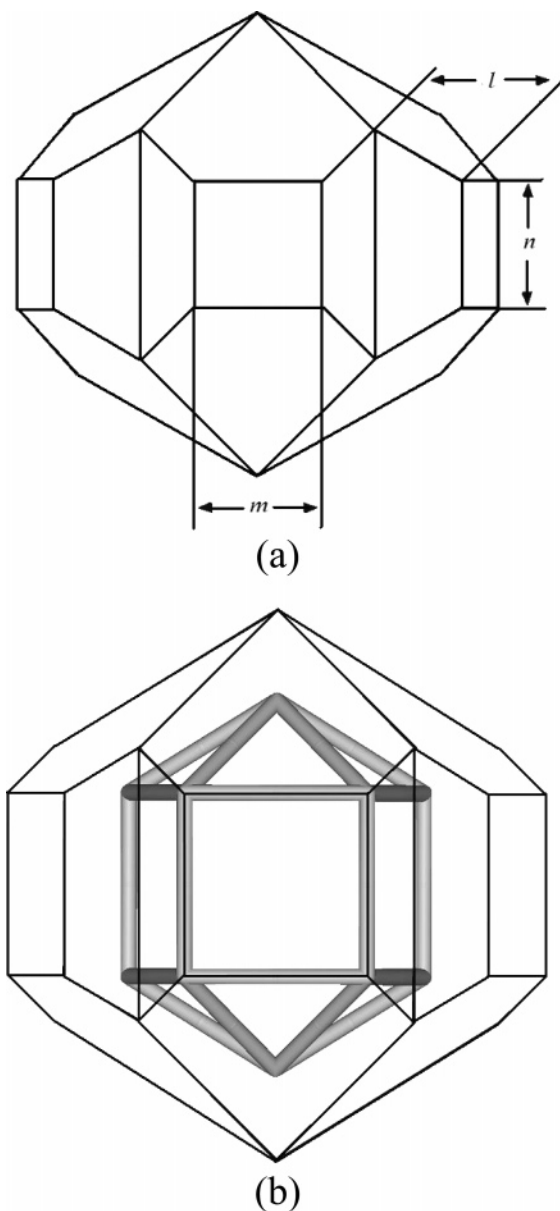


Figure 1. Configuration of (a) Ino decahedron and (b) m -Dh. m and n are width and height of the rectangular (100) facets, respectively, while l is the depth of the Marks re-entrance.

Dh, but there are Ih and fcc structures, which is consistent with the explanation of Baletto et al.¹² Together with the previous results of Ag_{13-120} , structural characteristic of Ag_{13-160} are investigated in detail for a better understanding of the changing rule of the structures with the cluster size. Results show that, with the growth of the cluster size, the dominant structural motif changes from amorphous (Ag_{13-48}) to Ih (Ag_{49-61}) and then to Dh (Ag_{62-160}). Furthermore, the clusters with Dh motif were classified into different kinds of m -Dh submotifs. The Dh structure of a given cluster was found to be a result of the competition among different m -Dh submotifs.

2. Method

The DLS method was proposed by combining the basic ideas of unbiased and biased methods.^{25,26} A DLS run starts from a randomly generated and locally minimized structure of a cluster, and then repetitively finds structures with lower energy by “lattice construction” and “lattice searching”. The former operation constructs all possible vacant sites around the starting

structure, and the latter operation finds structures with lower energy by moving the atoms with higher energy to the vacant sites with lower energy. The merit of the DLS method is that, by using dynamic lattice strategy, the searching space is greatly reduced. However, with the increase of the cluster size, the searching space will increase too. For optimization of large clusters, further reduction of the searching space is still needed. With a simple consideration that the searching space can be reduced by reducing the number of atoms involved in the optimization, DLSc was proposed.²⁷ In the DLSc method, instead of generating the starting structure randomly, an inner core is used to generate the starting structure. Only the atoms in outer layers are involved in lattice searching. Thus, the searching space is further reduced compared with that in the previous DLS.

For silver clusters, there are mainly three kinds of ordered structural motifs, i.e., Ih, Dh, and fcc. So, in this study, Ih, Dh, and fcc inner cores are used, respectively, in the DLSc method. The Ih core was obtained by constructing the Ih lattices with the method in ref 27. For m -Dh motif, its structure can be described with three parameters, m , n , l , where, as shown in Figure 1a, m and n are the width and the height of the rectangular (100) facets, respectively, and l is the depth of the Marks reentrance.^{28,29} Therefore, an m -Dh structure can be described by (m, n, l) . At the same time, an Ino Dh can be taken as an inner core of an m -Dh, because it can be obtained by moving out the atoms in the outer shells from an m -Dh (as shown in Figure 1b). Therefore, in this work, Ino Dh structure was used as a kind of core in DLSc for m -Dh motifs, and it was constructed by the method in ref 30. Furthermore, because fcc structures, including the truncated octahedron (TO), were found in silver clusters, fcc octahedron was also used as a kind of core, and it was constructed by the method in ref 30. Therefore, Ino Dh, Ih, and fcc octahedron were used as the inner cores of the starting structure in DLSc. The final result was obtained by choosing the structure with the lowest energy among the independent runs with different cores.

On the other hand, different semiempirical potentials have been proposed for transition and noble metallic clusters, such as effective-medium theory,³¹ the glue model,³² embedded-atom,³³ Sutton–Chen,³⁴ Rosato–Guillope–Legrand (RGL),³⁵ and the Gupta potential.³⁶ The later one was used in this work. Gupta potential is based on the second moment approximation of the electron density of states in the tight-binding (TB) theory. It can be depicted in the following form³⁶

$$V = \frac{U_N}{2} \sum_{i=1}^N V_i \quad (1)$$

where N is the number of atoms in the cluster and U_N is a function of the atom number N . V_i consists of a pairwise repulsion energy of Born–Mayer type and a N -body attractive contribution

$$V_i = A \sum_{j \neq i} \exp \left[-p \left(\frac{r_{ij}}{r_0} - 1 \right) \right] - \left(\sum_{j \neq i} \exp \left[-2q \left(\frac{r_{ij}}{r_0} - 1 \right) \right] \right)^{1/2} \quad (2)$$

r_{ij} is the distance between atom i and j , and r_0 is the equilibrium nearest-neighbor distance in the bulk metal. The parameters p and q represent the repulsive and the attractive interaction range, respectively. The parameter A is fitted to experimental values

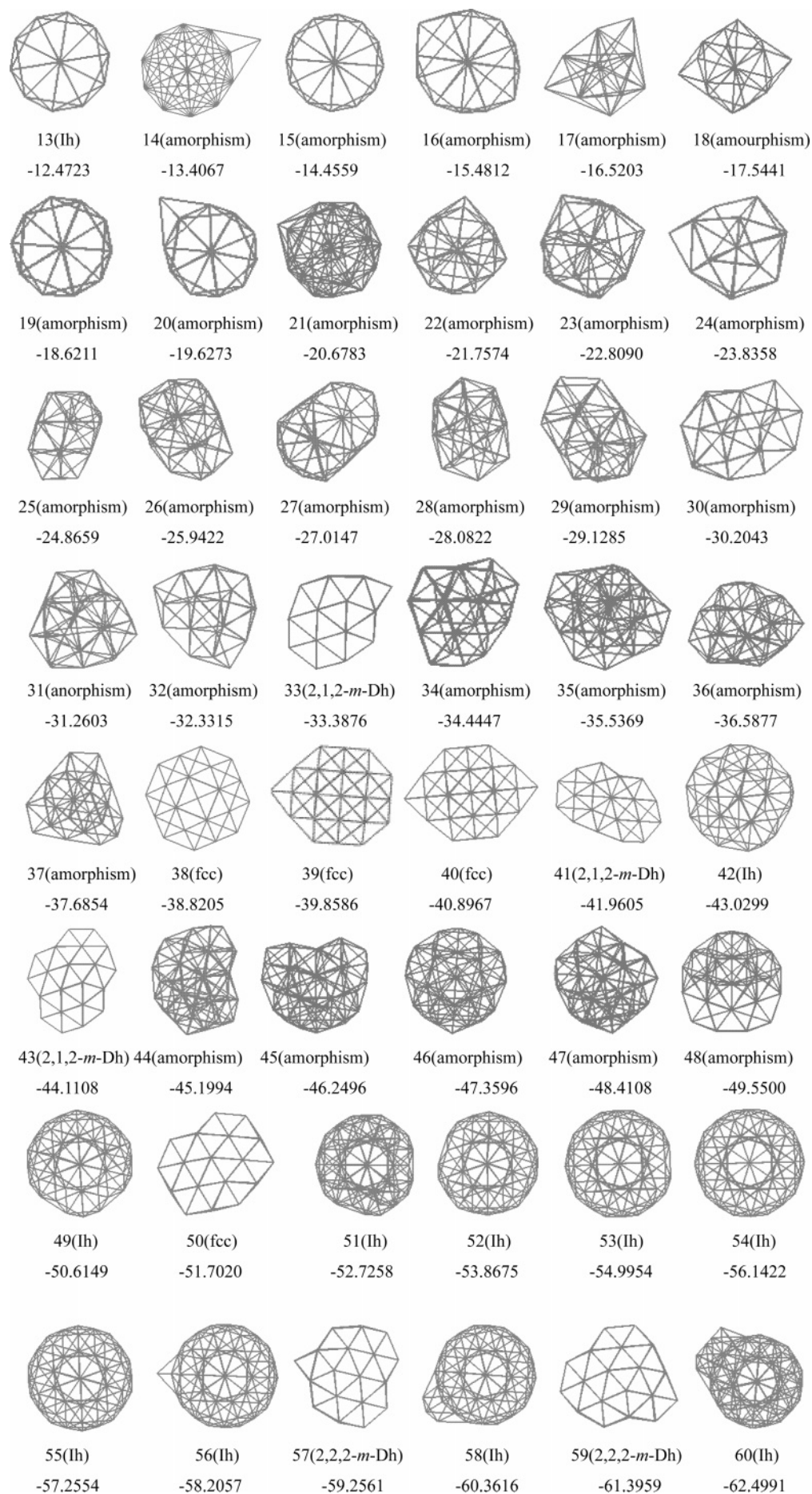


Figure 2. Part 1 of 4.

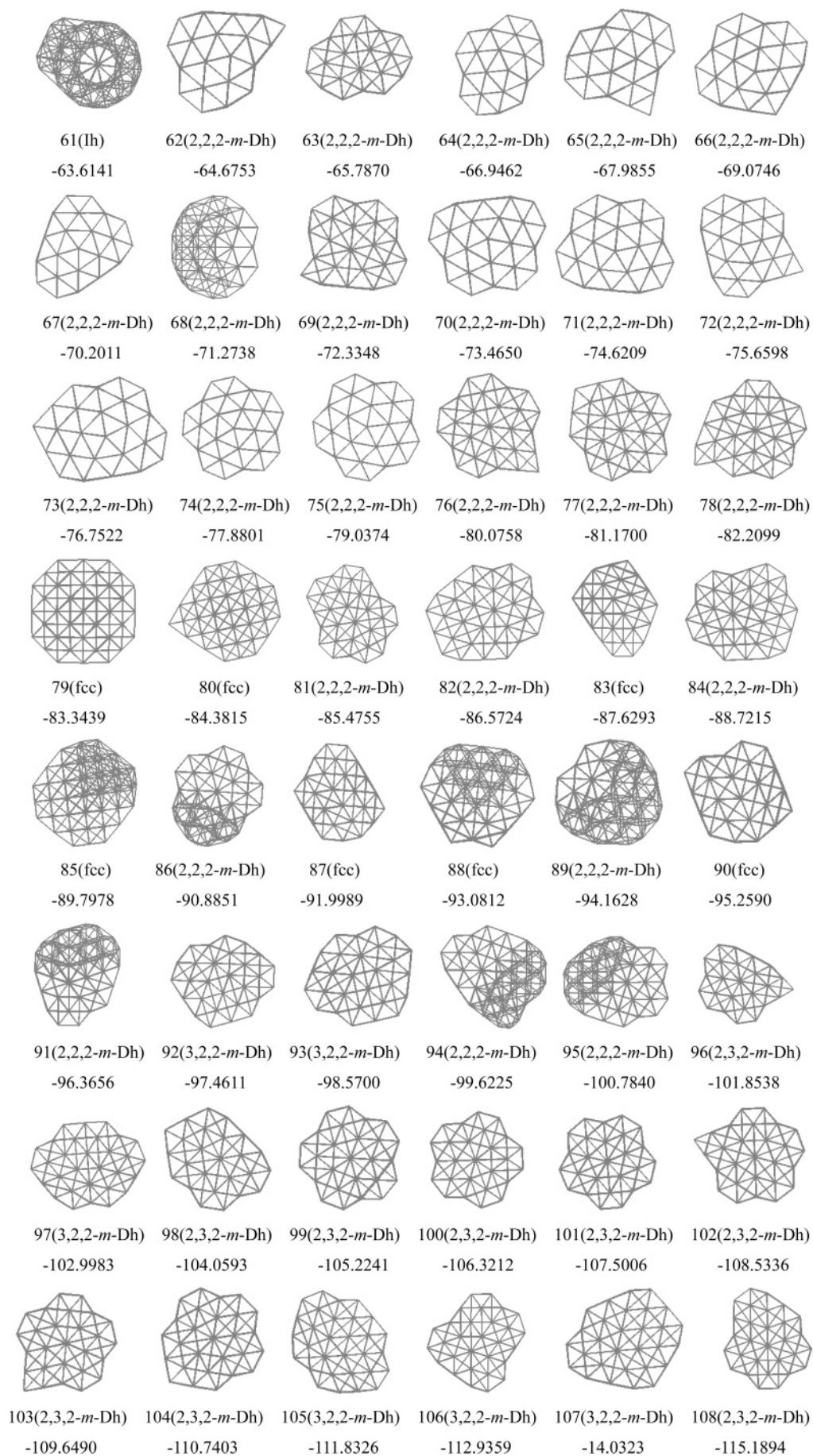


Figure 2. Part 2 of 4.

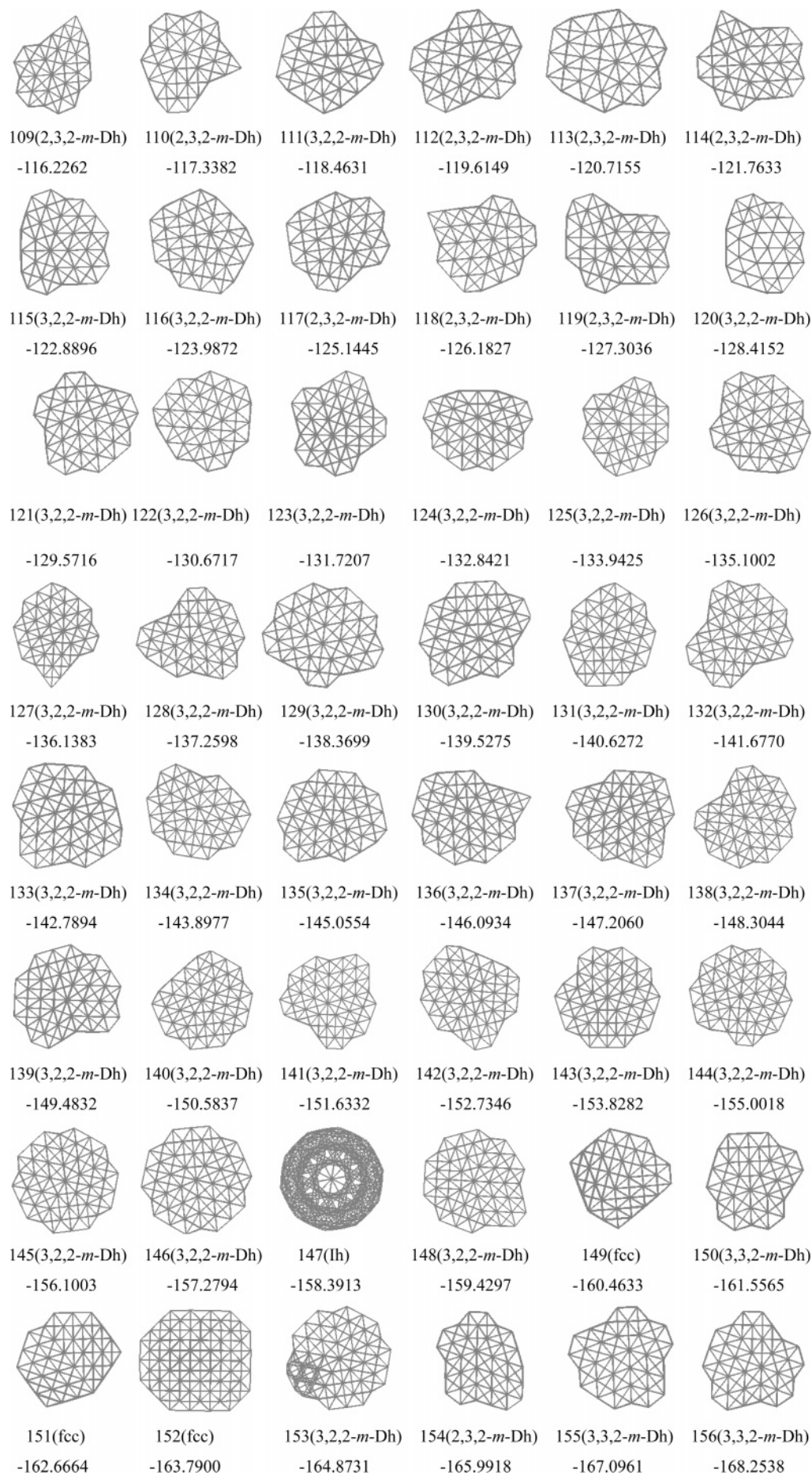


Figure 2. Part 3 of 4.

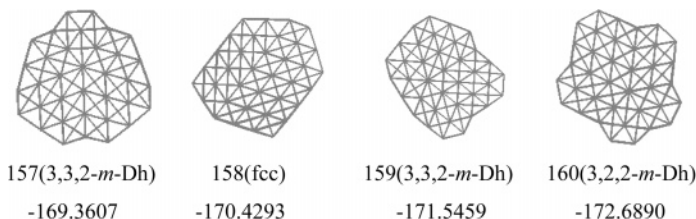


Figure 2. Part 4 of 4. Optimized structures, structural motifs, and the corresponding potential energies of the silver clusters from Ag₁₃ to Ag₁₆₀.

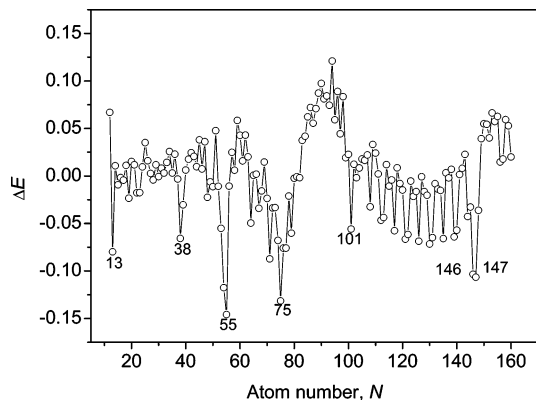


Figure 3. The finite difference (ΔE) of the energy of silver clusters from Ag₁₃ to Ag₁₆₀ for $E_{J(N)} = -0.92253 + 0.81546N^{1/3} + 0.29672N^{2/3} - 1.15599N$.

of the cohesive energy. In this work, $A = 0.09944$, $p = 10.12$, $q = 3.37$, and the reduced units with $r_0 = 1$ and $U_N = 1$ were used.³⁴

3. Results and Discussion

3.1. Putative Stable Structures of the Silver Clusters up to 160 Atoms. Figure 2 shows the optimized structures and the corresponding lowest energies of silver clusters from Ag₁₃ to Ag₁₆₀. Among these structures, Ag₁₃–Ag₁₂₀ were obtained with RTA and DLS methods, respectively, in our previous work,^{23,24} and Ag₁₂₁–Ag₁₆₀ were obtained with DLSc method in this work. It can be seen that, in the size range of 121–160, Dh is the dominant motif; only one Ih structure is obtained at Ag₁₄₇ and four fcc structures are obtained at the size of $N = 149, 151, 152,$ and 158 . This result is consistent with the results of the experiment and theoretical simulation.^{9,11} In the experiment of ref 9, the silver clusters are found to be mainly *m*-Dh in the range of 100–200, and, in the theoretical study of ref 11, it was found that the silver clusters up to 150 atoms are most possibly *m*-Dh.

3.2. Variation of the Structures and Energies with Size. Figures 3 and 4 shows the finite difference of energy ΔE and the second finite difference of energy $\Delta_2 E$ as a function of

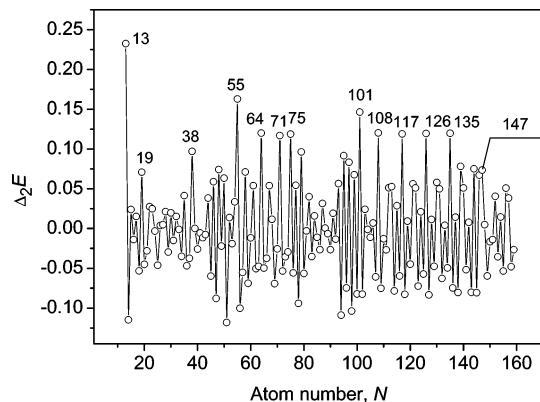


Figure 4. The second finite difference ($\Delta_2 E$) of the energy of silver clusters from Ag₁₃ to Ag₁₆₀.

cluster size N , respectively. The ΔE and $\Delta_2 E$ have the form as follows

$$\Delta E_{(N)} = E_{(N)} - E_{J(N)} \quad (3)$$

$$\Delta_2 E_{(N)} = E_{(N+1)} + E_{(N-1)} - 2E_{(N)} \quad (4)$$

where $E_{J(N)} = a + bN^{1/3} + cN^{2/3} + dN$ is a four-parameter fit of the energy of global minimum. Clearly, ΔE shows the variation of the cluster stability with cluster size, and $\Delta_2 E$ measures the stability of an N -atom cluster structure with respect to its neighboring cluster size. The negative peaks (or valleys) in Figure 3 and the positive peaks in Figure 4 indicate particularly stable structures comparing to their neighbors.

At first, in Figure 3, seven apparent valleys at Ag₁₃, Ag₃₈, Ag₅₅, Ag₇₅, Ag₁₀₁, Ag₁₄₆, and Ag₁₄₇ can be clearly found, which correspond to the complete structural motifs or the magic number clusters. Figure 5 shows the top view and side view of these structures. In more detail, from Ag₁₄ to Ag₄₈, the value of ΔE goes up very slightly with fluctuation except for the valley at Ag₃₈. The dominant motif in this size range is disordered structure, but the structural motif is fcc for Ag₃₈, Ag₃₉, and Ag₄₀. Another negative peak in Figure 3 can be found around Ag₅₅ from Ag₄₉ to Ag₆₁. In this range, the dominant motif is Ih, and

TABLE 1: Structural Distribution of Ag₁₃–160 Clusters

structural motif	n^a	cluster size (N)
amo rphi sm	28	14–32, 34–37, 44–48
Ih	13	13, 42, 49, 51–56, 58, 60, 61, 147
fcc	15	38–40, 50, 79, 80, 83, 85, 87, 88, 90, 149, 151, 152, 158
Dh		
(2,1,2) <i>m</i> -Dh	3	33, 41, 43
(2,2,2) <i>m</i> -Dh	27	57, 59, 62–78, 81, 82, 84, 86, 89, 91, 94, 95
(2,3,2) <i>m</i> -Dh	18	96, 98–104, 108–110, 112–114, 117–119, 154
(3,2,2) <i>m</i> -Dh	39	92, 93, 97, 105–107, 111, 115, 116, 120–146, 148, 153, 160
(3,3,2) <i>m</i> -Dh	5	150, 155–157, 159

^a Total number of the motif.

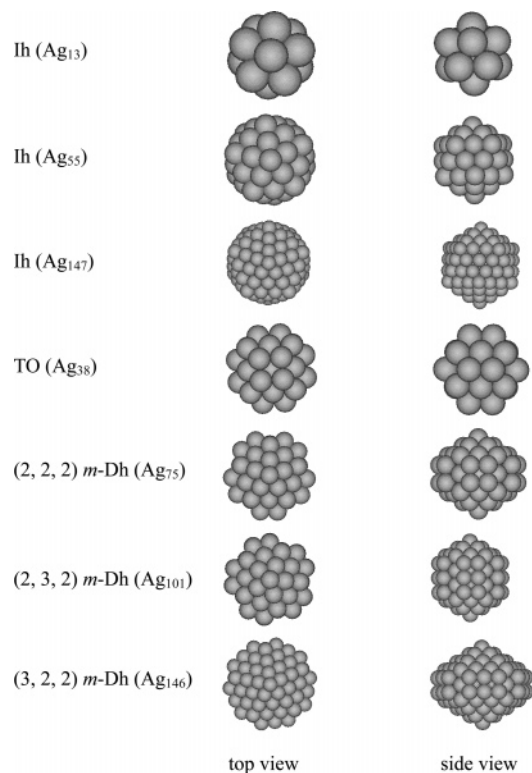


Figure 5. The top and side views of the silver cluster structures at magic numbers.

Ag_{55} is a complete Ih structure. However, the structural motif is fcc for Ag_{50} and (2,2,2) *m*-Dh for Ag_{57} and Ag_{59} . The clusters from Ag_{51} to Ag_{54} are more and more stable because that the structures gradually approach to the complete Ih structure, but the symmetry of Ag_{56} is destroyed by one extra atom in the out shell compared with the Ag_{55} complete Ih structure. Therefore, the negative peak is a reflection of the growth of Ih structure. When $N \geq 62$, the ΔE goes down with fluctuation and then goes up after $N = 75$ until $N = 95$. In this size range, the domain motif is (2, 2, 2) *m*-Dh and Ag_{75} is a complete one. After $N = 96$, there are a large smooth valley with comparatively bigger fluctuation and two sharp negative peaks at $N = 101$ and $146-147$, respectively. The large smooth valley indicates that there is no structural transition with the growth of the clusters in this size range, and the two sharp peaks indicate that there are complete structures at these corresponding sizes. It can be seen from Figure 2 that the dominant motif is *m*-Dh, but the structural motif is fcc for Ag_{149} , Ag_{151} , Ag_{152} , and Ag_{158} . In Figure 4, positive peaks corresponding to the negative peak in Figure 3 can be clearly found, e.g., at Ag_{13} , Ag_{38} , Ag_{55} , Ag_{75} , and Ag_{101} . However, there are some other peaks at Ag_{19} , Ag_{64} , Ag_{71} , Ag_{108} , Ag_{117} , Ag_{126} , and Ag_{135} in the figure, which show that these clusters are stable ones compared with their neighboring clusters. With a detail examination, it can be found that the structures of Ag_{19} , Ag_{64} , Ag_{71} , Ag_{126} , and Ag_{135} are more symmetrical than their neighboring clusters, e.g., Ag_{19} with D_{5h} symmetry but Ag_{18} with C_2 symmetry and Ag_{20} with C_{2v} symmetry. Ag_{108} and Ag_{117} have the more stable structures than their neighboring clusters due to the structures changing from (3,2,2) *m*-Dh to (2,3,2) *m*-Dh.

However, the motif of Ih with a single central vacancy was indicated being stable at Ag_{146} and Ag_{147} in ref 37. To further investigate the results, the Ih core with a single central vacancy was used to optimize the structures of Ag_{146} and Ag_{147} . It was found that the same structure as above are obtained. On the

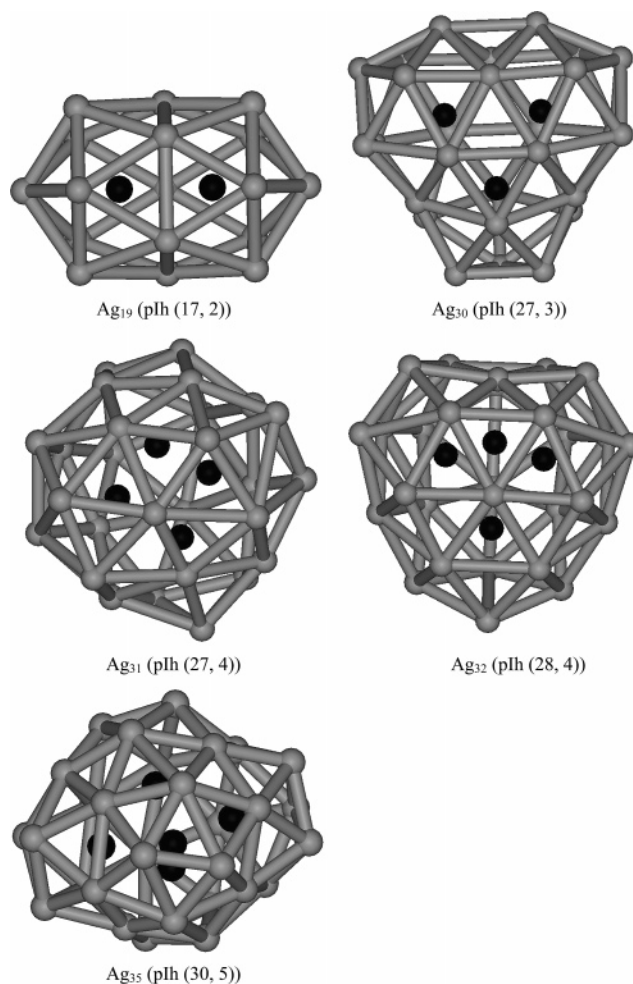


Figure 6. The structures of polyicosahedra motif. The atoms in the inner shell are shown in black.

other hand, in ref 11, the *m*-Dh motif has magic number at $N = 100$ with the corresponding (3,1,2) *m*-Dh structure. But there is no valley at $N = 100$ in Figure 3, and the putative stable structure is found to be (2,3,2) *m*-Dh with the potential energy -106.3212 for Ag_{100} in this work. For comparison, the energy of (3,1,2) *m*-Dh structure was calculated with Gupta potential and the result is -98.9675 . Therefore, (2,3,2) *m*-Dh should be more reliable structure for Ag_{100} .

3.3. Distribution of the Putative Stable Structures. The distribution of the Ag_{13-160} cluster structures is summarized in Table 1, including 28 amorphous structures, 13 Ih structures, 15 fcc structures, and 92 *m*-Dh structures in different sub-motifs.

Amorphous structures are obtained at the sizes of $N = 14-32$, $34-37$, and $44-48$. From Figure 2, it can be seen that these clusters have no symmetric axes. Among them, the structure of Ag_{15} is a complete icositetrahedron. It is classified into the amorphous motif for its scarceness. Furthermore, from Figure 2, it can also be found that the polyicosahedra (pIh)³⁸ motif can be obtained at Ag_{19} , Ag_{30-32} , and Ag_{35} . The structure of Ag_{19} has 30 facets, which is connected by two Ih structures with two atoms in the inner shell and 17 atoms in the outer shell, i.e., pIh (17, 2) as denoted in ref 38, and Ag_{20} has one more atom in outer shell than Ag_{19} . The structure of Ag_{30} is connected by three Ih, i.e., pIh (27, 3), the structures of Ag_{31} and Ag_{32} are connected by four Ih, i.e., pIh (27, 4) and pIh (28, 4), and the structure of Ag_{35} is connected by five Ih structures, i.e., pIh (30, 5). In order to more clearly demonstrate the structures of polyicosahedra motif, the structures of Ag_{19} , Ag_{20} ,

Ag₂₈, Ag_{30–32}, and Ag₃₅ were shown in Figure 6. In Figure 6, bonds among the atoms in inner shell and between the inner shell and outer shell were not drawn.

Ih motif clusters are obtained at sizes of $N = 13, 42, 49, 51–56, 58, 60, 61,$ and 147 . Obviously, most of the clusters with Ih motif are around the magic number 55. The structure of Ag₄₂ is a partial 55-atom Ih. The structures of Ag₅₁, Ag₅₂, and Ag₅₃ are constructed by removing one, two, and three atoms from the out shell of Ag₅₅, and on the contrary, that of Ag₅₆, Ag₅₈, Ag₆₀, and Ag₆₁ by adding one, three, five, and six atoms to the out shell, respectively. In Figure 2, Ag₅₄ and Ag₅₅ have the same top view, because Ag₅₄ is constructed by only removing the center atom from Ag₅₅.

Including the TO motif, 15 structures with the fcc motif are obtained at 38–40, 50, 79, 80, 83, 85, 87, 88, 90, 149, 151, 152, and 158. The structures of Ag₃₈ and Ag₇₉ are complete TO. Ag₃₉, Ag₄₀, Ag₈₀, and Ag₈₅ have the same structural motif because these structure are constructed with one and two extra atoms to the out shell of Ag₃₈, and one and six extra atoms to the out shell of Ag₇₉, respectively. Furthermore, it can be seen that the atomic arrangement in the outer shell of Ag₈₅ is different to the truncated octahedral motif.

From Table 1, it is obvious that the dominant motif of the clusters are m -Dh, especially when the size N is bigger than 57, and these m -Dh motif can be further divided into four sub-motifs of (2,2,2), (2,3,2), (3,2,2), and (3,3,2) m -Dh. The complete structures with these sub-motifs are found to be Ag₇₅, Ag₁₀₁, Ag₁₄₆, and Ag₁₉₂, respectively. In the small size range, only three clusters with (2,1,2) m -Dh motif at Ag₃₃, Ag₄₁, and Ag₄₃. The structures of Ag₃₃, Ag₄₁, and Ag₄₃ are partial (2,1,2) m -Dh. Furthermore, from Table 1 and Figure 2, it can be seen that the structures with fcc and (2,2,2) m -Dh motif are alternately obtained in the range of Ag₇₈ to Ag₉₁. Because 79 is a magic number of TO motif, fcc is the preferable motif for clusters around Ag₇₉. Then the preferable motif changes to (3,2,2) m -Dh at Ag₉₂, and back to (2,2,2) m -Dh at Ag₉₃. After Ag₉₅, the stable motif changes to (2, 3, 2) m -Dh in the range of Ag₉₆–Ag₁₁₉. So, from Ag₉₆ to Ag₁₁₉, the stable structure is a result of the competition between (2,3,2) m -Dh and (3,2,2) m -Dh. After Ag₁₁₉, the dominant motif changes to (3,2,2) m -Dh. With the increase of the cluster size, four (3,3,2) m -Dh structures are obtained at the sizes of $N = 150, 155, 157,$ and 159 . Therefore, stable structures with m -Dh motif are a result of the competition among the sub- m -Dh motifs. Because the magic number of (3,3,2) m -Dh is 192, it should be predictable that this motif will be a dominant motif for the clusters of $N \geq 150$.

4. Conclusion

The structures of silver clusters from Ag₁₂₁ to Ag₁₆₀ were optimized with DLSc, which is a variation of the previous DLS method. The new optimized results of Ag_{121–160} and our previous results of Ag_{13–120} were merged together to investigate the growth rule of silver clusters. A set of amorphous structures are obtained in the range of 13–48, together with several ordered structures at $N = 13, 15, 33, 38–40, 41,$ and 43 . In the size range of 49–61, the putative stable motif is Ih. The transition from Ih to Dh occurred in the size range of 57–62. Then, the putative stable motif change to m -Dh. On the hand, the stable structures of silver clusters with m -Dh motifs are found to be a result of the competition among the sub- m -Dh motifs.

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